



ON THE DETERMINANT OF AN HODGE-DE RHAM LIKE OPERATOR

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Abstract

Let (M, g) be a closed (i.e. compact without boundary), orientable, n -dimensional smooth Riemannian manifold, $C^\infty(M)$ the real algebra of smooth real functions on M , $A^k(M)$ the $C^\infty(M)$ -module of smooth differential k -forms, $0 \leq k \leq n$, and \mathbf{h} a pointwise nonsingular tensor field of type $(1,1)$ on M with vanishing Nijenhuis tensor. For such \mathbf{h} there is an associated exterior derivative $d_{\mathbf{h}}^{(k)} : A^k(M) \rightarrow A^{k+1}(M)$ having an adjoint $\delta_{\mathbf{h}}^{(k+1)} : A^{k+1}(M) \rightarrow A^k(M)$ with respect to the usual global inner product, so that one can define a (strongly) elliptic self-adjoint second order differential operator $\Delta_{\mathbf{h}}^{(k)} : A^k(M) \rightarrow A^k(M)$, which is a generalization of the Hodge-de Rham operator (see [4]). In this paper we shall discuss the smooth dependence on the Riemannian metric of the determinant of this \mathbf{h} -dependent Hodge-de Rham operator.

1 Introduction

Let (M, g) be a closed (i.e. connected, compact, and without boundary), oriented, n -dimensional smooth Riemannian manifold, $C^\infty(M)$ the real algebra of smooth real functions on M and $A^k(M)$ the $C^\infty(M)$ -module of smooth differential k -forms, $0 \leq k \leq n$. A tensor field \mathbf{h} of type $(1, 1)$ on M , which can be conceived as a $C^\infty(M)$ -linear mapping $\mathbf{h} : A^1(M) \rightarrow A^1(M)$, induces $C^\infty(M)$ -linear mappings $\mathbf{h}^{(q)} : A^k(M) \rightarrow A^k(M)$ for any nonnegative integer q , and the $\mathbf{h}^{(q)}$ are defined by setting $\mathbf{h}^{(q)} := 0$ if $q > k$, and

$$\mathbf{h}^{(q)}(\omega^1 \wedge \dots \wedge \omega^k) :=$$

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$$= \frac{1}{(k-q)!q!} \sum_{\sigma \in \mathfrak{S}_k} \operatorname{sgn}(\sigma) [\mathbf{h}(\omega^{\sigma(1)}) \wedge \dots \wedge \mathbf{h}(\omega^{\sigma(q)})] \wedge \omega^{\sigma(q+1)} \wedge \dots \wedge \omega^{\sigma(k)}$$

if $0 \leq q \leq k$, where $\omega^i \in A^1(M)$, $1 \leq i \leq k$, σ runs through all permutations of $\{1, \dots, k\}$, and $\operatorname{sgn}(\sigma)$ denotes the sign of the permutation σ . The transformation $\mathbf{h}^{(0)}$ is taken to be the identity mapping on $A^k(M)$.

If \mathbf{h} is a pointwise nonsingular tensor field of type $(1, 1)$ on M with vanishing Nijenhuis derivation, an \mathbf{h} -dependent exterior derivation $d_{\mathbf{h}}^{(k)} : A^k(M) \rightarrow A^{k+1}(M)$ can be defined by setting $d_{\mathbf{h}}^{(k)} := \mathbf{h}^{(1)} \circ d^{(k)} - d^{(k)} \circ \mathbf{h}^{(1)}$, where $d^{(k)} : A^k(M) \rightarrow A^{k+1}(M)$ is the usual exterior differential. When \mathbf{h} is the identity, $d_{\mathbf{h}}^{(k)}$ coincides with $d^{(k)}$. Let $\delta^{(k+1)} : A^{k+1}(M) \rightarrow A^k(M)$ be the adjoint of $d^{(k)}$ relative to the usual inner product induced by g . If $\mathbf{h}_t^{(1)}$ denotes the transpose of $\mathbf{h}^{(1)}$, the mapping $\delta_{\mathbf{h}}^{(k+1)} : A^{k+1}(M) \rightarrow A^k(M)$, defined by setting $\delta_{\mathbf{h}}^{(k+1)} := \delta^{(k+1)} \circ \mathbf{h}_t^{(1)} - \mathbf{h}_t^{(1)} \circ \delta^{(k+1)}$, is the adjoint of $d_{\mathbf{h}}^{(k)}$ with respect to the usual inner product. The \mathbf{h} -dependent Hodge-de Rham operator $\Delta_{\mathbf{h}}^{(k)} : A^k(M) \rightarrow A^k(M)$, defined by setting $\Delta_{\mathbf{h}}^{(k)} := d_{\mathbf{h}}^{(k-1)} \circ \delta_{\mathbf{h}}^{(k)} + \delta_{\mathbf{h}}^{(k+1)} \circ d_{\mathbf{h}}^{(k)}$, is elliptic and self-adjoint, and it reduces to the usual Hodge-de Rham operator $\Delta^{(k)}$ in the case when \mathbf{h} is the identity (see [4]).

2 Spectral properties of the \mathbf{h} -dependent Hodge-de Rham operators

In what follows we assume that \mathbf{h} is a nonsingular tensor field of type $(1, 1)$ on M with vanishing Nijenhuis derivation. Since $\Delta_{\mathbf{h}}^{(k)} : A^k(M) \rightarrow A^k(M)$ is an elliptic second order differential operator, formally self-adjoint and formally positive, and (M, g) is a closed, smooth Riemannian manifold, the following statements are valid (see also [2]).

Theorem 2.1 (i) For each $k \in \{0, 1, \dots, n\}$, there exists a discrete spectral resolution $\{\omega_{\mathbf{h};j}^{(k)}, \lambda_{\mathbf{h};j}^{(k)}\}_{j \in \mathbb{N}}$ of the operator $\Delta_{\mathbf{h}}^{(k)} : A^k(M) \rightarrow A^k(M)$, where $\omega_{\mathbf{h};j}^{(k)} \in A^k(M)$ for any $j \in \mathbb{N}$, that is $\{\omega_{\mathbf{h};j}^{(k)}\}_{j \in \mathbb{N}}$ is a complete orthonormal system in the real Hilbert space $L^2(A^k(M)) = W^{0,2}(A^k(M))$ such that $\Delta_{\mathbf{h}}^{(k)} \omega_{\mathbf{h};j}^{(k)} = \lambda_{\mathbf{h};j}^{(k)} \omega_{\mathbf{h};j}^{(k)}$ for any $j \in \mathbb{N}$. Moreover, $\lambda_{\mathbf{h};j}^{(k)} \in [0, +\infty)$ for any $j \in \mathbb{N}$, each eigenspace of $\Delta_{\mathbf{h}}^{(k)}$ is finite dimensional, and $0 \in \mathbb{R}$ is an eigenvalue of $\Delta_{\mathbf{h}}^{(k)} : A^k(M) \rightarrow A^k(M)$ if and only if the k -th Betti number $\beta_k(M) \neq 0$;

(ii) For any $m \in \mathbb{N}$, there exists $\ell(m) \in \mathbb{N}^*$ and a constant $C_m > 0$, such that

$$\|\omega_{\mathbf{h};j}^{(k)}\|_{\infty, m} \leq C_m (1 + (\lambda_{\mathbf{h};j}^{(k)})^{\ell(m)}),$$

where

$$\|\omega_{\mathbf{h};j}^{(k)}\|_{\infty,m} := \sup_{x \in M} |j_m(\omega_{\mathbf{h};j}^{(k)})_x|_{J^m(\Lambda^k T^* M)_x},$$

$j_m(\omega_{\mathbf{h};j}^{(k)})_x$ denoting the m -jet of the smooth section $\omega_{\mathbf{h};j} \in C^\infty(\Lambda^k T^* M)$ at the point $x \in M$;

(iii) If one arrange the eigenvalues of $\Delta_{\mathbf{h}}^{(k)}$ such that

$$0 \leq \lambda_{\mathbf{h};0}^{(k)} \leq \lambda_{\mathbf{h};1}^{(k)} \leq \dots,$$

then there exist real constants $C(k) > 0$ and $\varepsilon(k) > 0$ such that $\lambda_{\mathbf{h};j}^{(k)} \geq C(k)j^{\varepsilon(k)}$ if $j \geq j_0$ is sufficiently large;

(iv) Let $a_j^{(k)} := \langle \theta, \omega_{\mathbf{h};j}^{(k)} \rangle \in \mathbb{R}$, $j \in \mathbb{N}$, be the Fourier coefficients associated to $\theta \in L^2(A^k(M))$. If $\theta \in A^k(M)$, then

$$\sum_{j \in \mathbb{N}} |a_j^{(k)}| \lambda_{\mathbf{h};j}^{(k)} < +\infty$$

and the series $\sum_{j \in \mathbb{N}} |a_j^{(k)}| \lambda_{\mathbf{h};j}^{(k)}$ tends to θ uniformly with respect to the norm $\|\cdot\|_{\infty,i}$ for any $k \in \{0, 1, \dots, n\}$.

For an initial smooth differential k -form $\theta \in A^k(M)$, let us consider the heat equation associated to $\Delta_{\mathbf{h}}^{(k)} : A^k(M) \rightarrow A^k(M)$ with the initial condition θ :

$$\left(\frac{\partial}{\partial t} + \Delta_{\mathbf{h}}^{(k)}\right)\omega(x, t) = 0, \quad (1)$$

$$\lim_{t \searrow 0} \omega(x, t) = \theta(x), \quad (2)$$

where $x \in M$, $t \in (0, +\infty)$. Let $(e^{-t\Delta_{\mathbf{h}}^{(k)}})(\theta)$ be the unique solution of the evolution equation (1) that satisfies the initial condition (2). The linear operator $e^{-t\Delta_{\mathbf{h}}^{(k)}} : A^k(M) \rightarrow A^k(M)$ can be extended to a compact, self-adjoint operator from $L^2(A^k(M))$ into $L^2(A^k(M))$, denoted also by $e^{-t\Delta_{\mathbf{h}}^{(k)}}$. Let $\{\omega_{\mathbf{h};j}^{(k)}, \lambda_{\mathbf{h};j}^{(k)}\}$ be a discrete spectral resolution of $\Delta_{\mathbf{h}}^{(k)} : A^k(M) \rightarrow A^k(M)$ and let us consider the Fourier expansion of $\theta \in A^k(M)$:

$$\theta = \sum_{j \in \mathbb{N}} a_j^{(k)} \omega_{\mathbf{h};j}^{(k)}, \quad \text{where } a_j^{(k)} := \langle \theta, \omega_{\mathbf{h};j}^{(k)} \rangle, \quad j \in \mathbb{N}.$$

Let us define $\omega_{\mathbf{h};j}^{*(k)} \in (A^k(M))^*$ by $\omega_{\mathbf{h};j}^{*(k)}(\varpi) := \langle \varpi, \omega_{\mathbf{h};j}^{(k)} \rangle$, $\varpi \in A^k(M)$, such that $\|\omega_{\mathbf{h};j}^{*(k)}\| = \|\omega_{\mathbf{h};j}^{(k)}\|$. With these notations, the following statement is valid.

Theorem 2.2 *The following two series converge uniformly in the C^ℓ -topology for any $\ell \in \mathbb{N}$ if $t \geq \delta > 0$:*

$$E_{\mathbf{h}}^{(k)}(x, y, t) := \sum_{j \in \mathbb{N}} e^{-t\lambda_{\mathbf{h};j}^{(k)}} \omega_{\mathbf{h};j}^{(k)}(x) \otimes \omega_{\mathbf{h};j}^{*(k)}(y) \in \text{Hom}(\Lambda^k(T_y^*M), \Lambda^k(T_x^*M))$$

and

$$(e^{-t\Delta_{\mathbf{h}}^{(k)}}(\theta))(x, t) := \sum_{j \in \mathbb{N}} e^{-t\lambda_{\mathbf{h};j}^{(k)}} a_j^{(k)} \omega_{\mathbf{h};j}^{(k)}(x) := \int_M (E_{\mathbf{h}}^{(k)}(x, y, t)) \theta(y) d\mu_g(y),$$

where $x, y \in M$, and μ_g denotes the canonical measure on M associated to g .

Let us define the matrix of Fourier coefficients of the bounded linear operator

$$e^{-t\Delta_{\mathbf{h}}^{(k)}} : L^2(A^k(M)) \rightarrow L^2(A^k(M))$$

by

$$\begin{aligned} a_{ij}(E_{\mathbf{h}}^{(k)}) &:= \langle e^{-t\Delta_{\mathbf{h}}^{(k)}}(\omega_{\mathbf{h};i}^{(k)}), \omega_{\mathbf{h};j}^{(k)} \rangle = \\ &= \int_{x \in M} \int_{y \in M} (E_{\mathbf{h}}^{(k)}(x, y, t)) (\omega_{\mathbf{h};i}^{(k)}(y) | \omega_{\mathbf{h};j}^{(k)}(x)) d\mu_g(x) d\mu_g(y), \end{aligned}$$

$i, j \in \mathbb{N}$, $k \in \{0, 1, \dots, n\}$, where $(|)$ denotes the pointwise inner product on $A^k(M)$ defined by using the fiber metric on $\Lambda^k T^*M$ induced by g .

Theorem 2.3 (i) *With the previous notations the following equalities*

$$\begin{aligned} \text{Trace}_{L^2}(e^{-t\Delta_{\mathbf{h}}^{(k)}}) &:= \int_{x \in M} \text{Trace}_{\Lambda^k(T_x^*M)}(E_{\mathbf{h}}^{(k)}(x, x, t)) d\mu_g(x) = \\ &= \sum_{j \in \mathbb{N}} a_{jj}(E_{\mathbf{h};j}^{(k)}) = \sum_{j \in \mathbb{N}} e^{-t\lambda_{\mathbf{h};j}^{(k)}} \end{aligned}$$

are valid for each $k \in \{0, 1, \dots, n\}$, that is the continuous linear operator $e^{-t\Delta_{\mathbf{h}}^{(k)}} : L^2(A^k(M)) \rightarrow L^2(A^k(M))$ is a trace class operator;

(ii) *If $(A(M), d_{\mathbf{h}})$ is the elliptic cochain complex previously defined, then*

$$\sum_{k=0}^n (-1)^k \text{Trace}_{L^2}(e^{-t\Delta_{\mathbf{h}}^{(k)}}) = \text{Index}(A(M), d_{\mathbf{h}}) = \chi(M),$$

where $\chi(M)$ denotes the Euler-Poincaré characteristic of M .

Finally, let us discuss the asymptotics of $\text{Trace}_{L^2}(e^{-t\Delta_{\mathbf{h}}^{(k)}})$ as $t \searrow 0$ and relate these asymptotics to $\text{Index}(A(M), d_{\mathbf{h}})$.

For each $t \in (0, 1)$ and $k \in \{0, 1, \dots, n\}$, the restriction of $E_{\mathbf{h}}^{(k)}(\cdot, \cdot, t)$ on the diagonal of $M \times M$ admits an asymptotic expansion. More precisely, the following statement is valid.

Lemma 2.4 *For each $k \in \{0, 1, \dots, n\}$ and each $m \in \mathbb{N}$, there exists a smooth mapping $e_m(\cdot, \Delta_{\mathbf{h}}^{(k)}) : M \rightarrow \text{End}(\Lambda^k(T^*M))$ such that*

- (i) $e_m(x, \Delta_{\mathbf{h}}^{(k)}) \in \text{End}(\Lambda^k(T_x^*M))$ depends functorially on a finite number of jets of the symbol of the second order differential operator $\Delta_{\mathbf{h}}^{(k)}$;
- (ii) $E_{\mathbf{h}}^{(k)}(x, x, t) \sim \sum_m e_m(x, \Delta_{\mathbf{h}}^{(k)}) t^{\frac{m-n}{2}}$ as $t \searrow 0$, for any $x \in M$;
- (iii) $e_m(x, \Delta_{\mathbf{h}}^{(k)}) = 0$ for any $x \in M$ and any m odd.

For the proof, we refer to [5], Lemma 1.8.2.

Let

$$M \ni x \mapsto a_m(x, \Delta_{\mathbf{h}}^{(k)}) := \text{Trace}(e_m(x, \Delta_{\mathbf{h}}^{(k)})) \in \mathbb{R}$$

and

$$a_m(\Delta_{\mathbf{h}}^{(k)}) := \int_{x \in M} a_m(x, \Delta_{\mathbf{h}}^{(k)}) d\mu_g(x) \in \mathbb{R}, \quad (3.3)$$

the invariant scalar functions and the numerical invariants respectively associated to the differential operator $\Delta_{\mathbf{h}}^{(k)}$, $k \in \{0, 1, \dots, n\}$.

Lemma 2.5 (i) *The invariants $e_m(x, \Delta_{\mathbf{h}}^{(k)})$ and $a_m(x, \Delta_{\mathbf{h}}^{(k)})$, m even, can be expressed by local formulas that are homogeneous of order m in the jets of the total symbol of $\Delta_{\mathbf{h}}^{(k)}$ for each $k \in \{0, 1, \dots, n\}$;*

- (ii) $a_m(x, \Delta_{\mathbf{h}}^{(k)}) = 0$ for each m odd and any $k \in \{0, 1, \dots, n\}$;
- (iii) $\text{Trace}_{L^2}(e^{-t\Delta_{\mathbf{h}}^{(k)}}) \sim \sum_m a_m(\Delta_{\mathbf{h}}^{(k)}) t^{\frac{m-n}{2}}$ as $t \searrow 0$, for each $k \in \{0, 1, \dots, n\}$.

For the proof of statement (i) we refer to [5], Lemma 1.8.3(c), while the statement (ii) is an immediate consequence of Lemma 2.4(iii). Statement (iii) is an immediate consequence of the definition of $\text{Trace}_{L^2}(e^{-t\Delta_{\mathbf{h}}^{(k)}})$ [see Theorem 2.3(i)] and of Lemma 2.4(ii).

3 The determinant of an h-dependent Hodge-de Rham operator

In this Section we assume that the k -th Betti number $\beta_k(M)$ of M is equal to zero, so we do not must to deal with the 0-spectrum. Let \mathcal{R} be a complementary region to a cone about positive real axis lying in the right half plane

\mathbb{C} , cone that includes the symbolic spectrum of $\Delta_{\mathbf{h}}^{(k)}$, so that the intersection of this symbolic spectrum with \mathcal{R} is empty. We normalize the choice of \mathcal{R} so that the boundary γ is a curve about the positive real axis which consists of a portion of a large circle around the origin and of two rays lying in the right half plane. We orient γ in a clockwise fashion. If $\lambda \in \mathcal{R}$ and $|\lambda|$ is large, then $(\Delta_{\mathbf{h}}^{(k)} - \lambda)$ is invertible. Since $\Delta_{\mathbf{h}}^{(k)}$ is positive definite, we may assume $Re(\lambda) > 0$ for $\lambda \in \gamma$. Choose the branch of λ^s , $s \in \mathbb{C}$, defined on the right half plane so $1^s = 1$. The L^2 -norm of the operator $(\Delta_{\mathbf{h}}^{(k)} - \lambda)^{-1}$ is uniformly bounded in $L^2(A^k(M))$. For $Re(s) > 1$, the integral

$$(\Delta_{\mathbf{h}}^{(k)})^{-s} := (2\pi i)^{-1} \int_{\gamma} \lambda^{-s} (\Delta_{\mathbf{h}}^{(k)} - \lambda)^{-1} d\lambda$$

converges to a bounded operator on $L^2(A^k(M))$.

Let $\{\omega_{\mathbf{h};j}^{(k)}, \lambda_{\mathbf{h};j}^{(k)}\}_{j \in \mathbb{N}}$ be a discrete spectral resolution of $\Delta_{\mathbf{h}}^{(k)}$ [see Theorem 2.1(i)], and

$$E_{\mathbf{h}}^{(k)}(x, y, s, \Delta_{\mathbf{h}}^{(k)}) := \sum_j (\lambda_{\mathbf{h};j}^{(k)})^{-s} \omega_{\mathbf{h};j}^{(k)}(x) \otimes \omega_{\mathbf{h};j}^{*(k)}(y) \quad (3)$$

the kernel of $(\Delta_{\mathbf{h}}^{(k)})^{-s}$. Theorem 2.1 implies the existence of real constants $\delta > 0$ and $\ell(m) > 0$ such that

$$\lambda_{\mathbf{h};j}^{(k)} \geq j^{\delta} \text{ and } \|\omega_{\mathbf{h};j}^{(k)}\|_{\infty, m} \leq (\lambda_{\mathbf{h};j}^{(k)})^{\ell(m)}, \quad (4)$$

whence it follows that (3) converges absolutely and uniformly in the $\|\cdot\|_{\infty, m}$ -norm, for $Re(s)$ large.

Definition 3.1 *We define the generalized zeta function*

$$\begin{aligned} \zeta(s, \Delta_{\mathbf{h}}^{(k)}) &:= \text{trace}_{L^2}((\Delta_{\mathbf{h}}^{(k)})^{-s}) = \sum_j (\lambda_{\mathbf{h};j}^{(k)})^{-s} = \\ &= \int_M \text{Trace}(E_{\mathbf{h}}^{(k)}(x, x, s, \Delta_{\mathbf{h}}^{(k)})) d\mu_g(x). \end{aligned}$$

Using relations (4), one can show that it converges uniformly and absolutely for $Re(s)$ large.

We use the Mellin transform to relate the zeta function to the heat kernel. From Lemma 2.5(iii), Theorem 2.3(i) and the identity

$$\int_0^{\infty} t^{s-1} e^{-\lambda_{\mathbf{h}}^{(k)} t} dt = (\lambda_{\mathbf{h}}^{(k)})^{-s} \Gamma(s),$$

where $\Gamma : \{s \in \mathbb{C} \mid Re(s) > 0\} \rightarrow \mathbb{C}$, $\Gamma(s) := \int_0^{\infty} t^{s-1} e^{-t} dt$, we obtain the following theorem.

Theorem 3.2 (i) If $\operatorname{Re}(s)$ is large, then

$$\zeta(s, \Delta_{\mathbf{h}}^{(k)}) = \Gamma(s)^{-1} \int_0^\infty t^{s-1} \operatorname{Trace}_{L^2}(e^{-t\Delta_{\mathbf{h}}^{(k)}}) dt;$$

(ii) $\Gamma(s)\zeta(s, \Delta_{\mathbf{h}}^{(k)})$ has a meromorphic extension to \mathbb{C} with isolated simple poles at $s = \frac{n-m}{2}$ for $m \in \mathbb{N}$ and

$$\operatorname{Res}_{s=\frac{n-m}{2}} \Gamma(s)\zeta(s, \Delta_{\mathbf{h}}^{(k)}) = a_m(\Delta_{\mathbf{h}}^{(k)});$$

(iii) Let $\Delta_{\mathbf{h}}^{(k)}(\epsilon)$ be a smooth one parameter family of such operators. Then $\zeta(s, \Delta_{\mathbf{h}}^{(k)}(\epsilon))$ is smooth with respect to (s, ϵ) away from the exceptional values at $s = \frac{n-m}{2}$.

Remark 3.3 Because by the determinant of the \mathbf{h} -dependent Hodge-de Rham operator $\Delta_{\mathbf{h}}^{(k)}$ we understand the product of all eigenvalues, considered with their multiplicity, that is

$$\det(\Delta_{\mathbf{h}}^{(k)}) = \prod_{j=1}^{\infty} \lambda_{\mathbf{h},j}^{(k)},$$

by Theorem 3.2 one can define a regularization of it, using the generalized zeta function, namely

$$\det_{\zeta}(\Delta_{\mathbf{h}}^{(k)}) = e^{\left(-\frac{d}{ds}\Big|_{s=0} \zeta(s, \Delta_{\mathbf{h}}^{(k)})\right)}. \quad (5)$$

Note that this is formally correct since if there were only a finite number of eigenvalues we would have

$$\frac{d}{ds} \left(\sum_j (\lambda_{\mathbf{h},j}^{(k)})^{-s} \right) = - \sum_j (\lambda_{\mathbf{h},j}^{(k)})^{-s} \log(\lambda_{\mathbf{h},j}^{(k)})$$

which takes the value $-\log(\prod_j \lambda_{\mathbf{h},j}^{(k)})$ at $s = 0$.

In what follows, we endow the set $\mathfrak{M}(M)$ of all smooth Riemannian metrics on M with the structure of a Fréchet manifold (see [1] and [7]). In the paper [6] it is proven the following result.

Corollary 3.4 If M is a closed, n -dimensional smooth manifold, and \mathbf{h} a nonsingular tensor field of type $(1, 1)$ on M with vanishing Nijenhuis derivation, then for each fixed $t \in (0, \infty)$, the real functions

$$\operatorname{Trace}_{L^2}(e^{-t\Delta_{\mathbf{h}}^{(k)}(\cdot)}) = \sum_{j=1}^{\infty} e^{\lambda_{\mathbf{h},j}^{(k)}(\cdot)t} : \mathfrak{M}(M) \rightarrow \mathbb{R}$$

are smooth with respect to the canonical Fréchet manifold structure considered on $\mathfrak{M}(M)$.

Therefore, by relation (5), by Theorem 3.2 and by Corollary 3.4 one obtains

Corollary 3.5 *If M is a closed, n -dimensional smooth manifold, and \mathbf{h} a nonsingular tensor field of type $(1, 1)$ on M with vanishing Nijenhuis derivation, then for each fixed $t \in (0, \infty)$, the real functions*

$$\det_{\zeta}(\Delta_{\mathbf{h}}^{(k)}(\cdot)) = e^{\left(-\frac{d}{ds}\Big|_{s=0} \zeta(s, \Delta_{\mathbf{h}}^{(k)}(\cdot))\right)} : \mathfrak{M}(M) \rightarrow \mathbb{R}$$

are smooth with respect to the canonical Fréchet manifold structure considered on $\mathfrak{M}(M)$.

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